

## Functional differential equations for the $q$ -Fourier transform of $q$ -Gaussians

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

2010 J. Phys. A: Math. Theor. 43 095202

(<http://iopscience.iop.org/1751-8121/43/9/095202>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 171.66.16.158

The article was downloaded on 03/06/2010 at 08:57

Please note that [terms and conditions apply](#).

# Functional differential equations for the $q$ -Fourier transform of $q$ -Gaussians

S Umarov<sup>1</sup> and S M Duarte Queirós<sup>2</sup>

<sup>1</sup> Department of Mathematics, Tufts University, Medford, MA, USA

<sup>2</sup> Unilever R&D Port Sunlight, Quarry Road East, Wirral, CH63 3JW, UK

E-mail: [sdqueiro@gmail.com](mailto:sdqueiro@gmail.com)

Received 8 September 2009, in final form 7 January 2010

Published 15 February 2010

Online at [stacks.iop.org/JPhysA/43/095202](http://stacks.iop.org/JPhysA/43/095202)

## Abstract

In this paper the question ‘is the  $q$ -Fourier transform of a  $q$ -Gaussian a  $q'$ -Gaussian (with some  $q'$ ) up to a constant factor?’ is studied for the whole range of  $q \in (-\infty, 3)$ . This question is connected with applicability of the  $q$ -Fourier transform in the study of limit processes in nonextensive statistical mechanics. Using the functional differential equation approach we prove that the answer is affirmative if and only if  $1 \leq q < 3$ , excluding two particular cases of  $q < 1$ , namely  $q = \frac{1}{2}$  and  $q = \frac{2}{3}$ . Complementarily, we discuss some applications of the  $q$ -Fourier transform to nonlinear partial differential equations such as the porous medium equation.

PACS numbers: 02.30.Sa, 02.30.Uu, 89.65.Gh, 89.75.Da

Mathematics Subject Classification: 35Q84, 37A50, 43A32, 65L03

## 1. Introduction

Approximately a century after Boltzmann’s seminal works which have turned into the cornerstones of statistical mechanics, Tsallis [1] introduced an entropic form aimed to accommodate the description of systems whose fundamental features may not be fitted in the Boltzmann–Gibbs formalism (see details in [2–4]). Tsallis’ entropic form, which is usually called the non-additive  $q$ -entropy, recovers the classic Boltzmann–Gibbs entropic form,  $S(f) = -\int f(x)\ln f(x) dx$  in the limit case  $q \rightarrow 1$ . Concomitantly, there is the nonextensive statistical mechanics formalism based on  $q$ -algebra and the  $q$ -Gaussian probability density function, which maximizes  $q$ -entropy under certain appropriate constraints (see [1, 5] and references therein). Recently, the  $q$ -Fourier transform [6] was introduced as a tool for the study of attractors of strongly correlated random variables in conjunction with the  $q$ -central limit theorem. The existence of such a theorem within nonextensive statistical mechanics was first conjectured in [7, 8]. In this paper we shed light on the question—whether the  $q$ -Fourier transform of a  $q$ -Gaussian is a  $q'$ -Gaussian, clarifying thereupon applicability of the  $q$ -Fourier

transform technique as a mathematical tool. A key to this matter is crucial because the  $q$ -Fourier transform is relevant to the study of limit distributions of strongly correlated random variables, as well as to solutions of partial differential equations with physical significance. Moreover, a positive answer implies validating a mapping relation of  $q$  onto  $q'$  obtained from the  $q$ -Fourier transform. This relation has been predominant for the establishment of other stable distributions, namely the  $(q, \alpha)$ -stable distributions [9]. We recall that, by definition, the  $q$ -Fourier transform of a nonnegative  $f \in L_1(R)$  is defined by the formula

$$F_q[f](\xi) = \int_{\text{supp } f} e_q^{ix\xi} \otimes_q f(x) dx, \quad \xi \in (-\infty, \infty), \tag{1}$$

where  $q < 3$ , the symbol  $\otimes_q$  stands for the  $q$ -product and

$$e_q^z = (1 + (1 - q)z)^{1/(1-q)}, \quad z \in C, \tag{2}$$

is a  $q$ -exponential, which is the usual exponential function  $e^z$  in the limit  $q \rightarrow 1$ , and defined for all  $z \in C \setminus z_0$ ,  $z_0 = -1/(1 - q)$ , with principal values along the cut  $(-\infty, z_0)$  if  $q \neq 1$  (see [6, 7] for details). The equality

$$e_q^{ix\xi} \otimes_q f(x) = f(x) e_q^{\frac{ix\xi}{|f(x)|^{1-q}}},$$

valid for all  $x \in \text{supp } f$  implies the following representation for the  $q$ -Fourier transform without usage of the  $q$ -product:

$$F_q[f](\xi) = \int_{\text{supp } f} f(x) e_q^{ix\xi|f(x)|^{q-1}} dx. \tag{3}$$

The paper is organized as follows: in section 2 we mention some properties of the  $q$ -Fourier transform. In section 3 we derive functional differential equations for the  $q$ -Fourier transform of  $q$ -Gaussians. Then, based on the results of this section, we show that the answer to the above question is affirmative for all  $1 \leq q < 3$ , and for two particular values of  $q < 1$ , namely for  $q = 1/2$  and  $q = 2/3$ . We also show that if  $q < 1$ , except two values mentioned above, the  $q$ -Fourier transform of a  $q$ -Gaussian is no longer a  $q'$ -Gaussian,  $\forall q' < 3$ . A relevant physical application of the  $q$ -Fourier transform and the functional differential equations studied in section 3 is addressed in section 4.

## 2. Preliminaries

Representation (3) for the  $q$ -Fourier transform implies the following proposition.

**Proposition 2.1.** For any constants  $a > 0$ ,  $b \neq 0$ ,

- (i)  $F_q[af(x)](\xi) = aF_q[f(x)]\left(\frac{\xi}{a^{1-q}}\right)$ ;
- (ii)  $F_q[f(bx)](\xi) = \frac{1}{b}F_q[f(x)]\left(\frac{\xi}{b}\right)$ .

Let  $\beta$  be a positive number. By definition, the function

$$G_q(\beta; x) = \frac{\sqrt{\beta}}{C_q} e_q^{-\beta x^2}, \tag{4}$$

with domains given below, is called a  $q$ -Gaussian density function:

- (i) if  $q < 1$ , then  $G_q(\beta; x)$  is defined on the compact set  $[-K_\beta, K_\beta]$ , where  $K_\beta = (\beta(1 - q))^{-1/2}$ ;
- (ii) if  $1 \leq q < 3$ , then  $G_q(\beta; x)$  is defined on the whole real axis  $\mathbb{R} = (-\infty, \infty)$ .

In expression (4)  $C_q$  is the normalizing constant, i.e.  $C_q = \int_{-\infty}^{\infty} e_q^{-x^2} dx$ , whose explicit form is given by (see, e.g., [6])

$$C_q = \begin{cases} \frac{2\sqrt{\pi} \Gamma(\frac{1}{1-q})}{(3-q)\sqrt{1-q} \Gamma(\frac{3-q}{2(1-q)})}, & -\infty < q < 1, \\ \sqrt{\pi}, & q = 1, \\ \frac{\sqrt{\pi} \Gamma(\frac{3-q}{2(q-1)})}{\sqrt{q-1} \Gamma(\frac{1}{q-1})}, & 1 < q < 3. \end{cases} \tag{5}$$

We use the convention  $K_\beta = \infty$  if  $q \geq 1$ , since, by definition, the support of the  $q$ -Gaussian is not bounded in this case.

Note that  $q$ -exponentials possess the property  $e_q^z \otimes_q e_q^w = e_q^{z+w}$  [12, 13]. This immediately implies the following assertion.

**Proposition 2.2.** For all  $q < 3$  the  $q$ -Fourier transform of  $e_q^{-\beta x^2}$ ,  $\beta > 0$ , can be written in the form

$$F_q[e_q^{-\beta x^2}](\xi) = \int_{-K_\beta}^{K_\beta} e_q^{-\beta x^2 + ix\xi} dx. \tag{6}$$

**Corollary 2.3.** Let  $q < 3$ . Then

$$F_q[e_q^{-\beta x^2}](\xi) = 2 \int_0^{K_\beta} e_q^{-\beta x^2} \cos_q\left(\frac{x\xi}{[e_q^{-\beta x^2}]^{1-q}}\right) dx, \quad \forall q,$$

where

$$\cos_q(x) = \frac{e_q^{ix} + e_q^{-ix}}{2}.$$

The assertion below was proved in [6].

**Proposition 2.4.** Let  $1 \leq q < 3$ . Then

$$F_q[G_q(\beta; x)](\xi) = e_{q_1}^{-\beta_* \xi^2}, \quad \xi \in \mathbb{R}, \tag{7}$$

where  $q_1 = \frac{1+q}{3-q}$  and  $\beta_* = \frac{3-q}{8\beta^{2-q} C_q^{2(q-1)}}$ .

**Proposition 2.5.** Let  $q < 1$ . Then

$$F_q[G_q(\beta, x)](\xi) = e_{q_1}^{-\beta_* \xi^2} \left(1 - \frac{2}{C_q} \operatorname{Im} \int_0^{d_\xi} e_q^{b_\xi + i\tau} d\tau\right), \quad \xi \in \left(-K_{\frac{1}{4\beta}}, K_{\frac{1}{4\beta}}\right),$$

where  $q_1$  and  $\beta_*$  are as in proposition 2.4 and  $b_\xi + id_\xi = \frac{K_\beta \sqrt{\beta} - i \frac{\xi}{2\sqrt{\beta}}}{\left[e_q^{-\frac{\xi^2}{4\beta}}\right]^{\frac{1-q}{2}}}$ .

**Proof.** The proof of this statement can be obtained applying the Cauchy theorem, that is by integrating the function  $e_q^{-\beta z^2 + iz\xi}$  over the closed contour  $C = C_0 \cup C_1 \cup C_- \cup C_+$ , where  $C_p = (-K_\beta + pi, K_\beta + ip)$ ,  $p = 0, 1$ , and  $C_\pm = [\pm K_\beta, \pm K_\beta + i]$ .  $\square$

Unifying propositions 2.4 and 2.5,

$$F_q[G_q(\beta, x)](\xi) = e_{q_1}^{-\beta_* \xi^2} + I_{(-\infty, 0)}(q) T_q(\xi),$$

where  $I_{(a,b)}(\cdot)$  designates the indicator function of an interval  $(a, b)$ , and

$$T_q(\xi) = -\frac{2}{C_q} e_{q_1}^{-\beta_* \xi^2} \operatorname{Im} \int_0^{d_\xi} e_q^{b_\xi + i\tau} d\tau.$$

Thus, for  $q \geq 1$ , the operator  $F_q$  transforms a  $q$ -Gaussian into a  $q_1$ -Gaussian with the factor  $C_{q_1} \beta^{-1/2}$ . However, for  $q < 1$  the additional tail  $T_q(\xi)$  appears.

Further, introduce a sequence  $q_n$  defined as

$$q_n = \frac{2q - n(q - 1)}{2 - n(q - 1)}, \tag{8}$$

where  $-\infty < n < \frac{2}{q-1} - 1$  if  $1 < q < 3$ , and  $n > -\frac{2}{1-q}$  if  $q < 1$ . Obviously,  $q_0 = q$ . Note also that if  $q = 1$ , then  $q_n = 1$  for all  $n = 0, \pm 1, \dots$ . Let  $\mathbb{Z}$  be the set of all integer numbers. Denote by  $\mathbb{N}_q$  a subset of  $\mathbb{Z}$  defined as

$$\mathbb{N}_q = \begin{cases} \{n \in \mathbb{Z} : n < \frac{2}{q-1} - 1\}, & \text{if } 1 < q < 3, \\ \mathbb{Z}, & \text{if } q = 1, \\ \{n \in \mathbb{Z} : n > -\frac{2}{1-q}\}, & \text{if } q < 1. \end{cases}$$

**Proposition 2.6.** For all  $n \in \mathbb{N}_q$  the following relations hold:

- (i)  $(3 - q_n)q_{n+1} = (3 - q_{n-2})q_n$ ,
- (ii)  $2C_{q_{n-2}} = \sqrt{q_n} (3 - q_n) C_{q_n}$ .

**Proof.**

(i) It follows from the definition of  $q_n$  that  $q_{n+1} = (1 + q_n)/(3 - q_n)$ . This yields

$$(3 - q_n)q_{n+1} = 1 + q_n = \left(1 + \frac{1}{q_n}\right) q_n. \tag{9}$$

Further, it is easy to verify that the equality  $q_{k-1} + q_{k+1}^{-1} = 2$  holds for all  $k \in \mathbb{N}_q$ . Applying this relationship for  $k = n - 1$ , we have  $1/q_n = 2 - q_{n-2}$ . Now taking this into account in (9), we obtain (i).

(ii) Obviously, for  $q = 1$  relationship (ii) reads  $2\sqrt{\pi} = 2\sqrt{\pi}$ . Let  $q \neq 1$ . Note that for any  $n \in \mathbb{N}_q$  the condition  $1 < q < 3$  implies  $1 < q_n < 3$ , as well as the condition  $q < 1$  implies  $q_n < 1$ . Using the explicit forms for  $C_q$  given in (5) and the relationship  $2 - q_{n-2} = 1/q_n$ , one obtains in the case  $1 < q < 3$

$$\frac{2C_{q_{n-2}}}{C_n} = \frac{\sqrt{q_n} \Gamma\left(\frac{1+q_n}{2(q_n-1)}\right)}{\frac{1}{2(q_n-1)} \Gamma\left(\frac{3-q_n}{2(q_n-1)}\right)} = \sqrt{q_n}(3 - q_n);$$

and in the case  $q < 1$

$$\frac{2C_{q_{n-2}}}{C_n} = \frac{\sqrt{q_n}(3 - q_n) \Gamma\left(\frac{3-q_n}{2(1-q_n)}\right)}{\frac{1+q_n}{2(1-q_n)} \Gamma\left(\frac{1+q_n}{2(1-q_n)}\right)} = \sqrt{q_n}(3 - q_n),$$

completing the proof of part (ii). □

### 3. Main results

#### 3.1. Functional differential equations

Let  $g_q(\beta, \xi)$  be the  $q$ -Fourier transform of a  $q$ -Gaussian  $G_q(\beta, \xi)$ , i.e.  $g_q(\beta, \xi) = F_q[G_q(\beta, x)](\xi)$ , and  $g_q(\xi) = g_q(1, \xi)$  for  $\beta = 1$ . Further, let  $Y_q(\xi) = F_q[e_q^{-x^2}](\xi)$ . In accordance with proposition 2.2,

$$Y_q(\xi) = \int_{-K}^K e_q^{-x^2+ix\xi} dx, \quad \xi \in \mathbb{R}, \tag{10}$$

where  $K = K_1 = \frac{1}{\sqrt{1-q}}$  if  $q < 1$ , and  $K = \infty$ , if  $q \geq 1$ .

**Lemma 3.1.** For any  $q < 3$  and  $\beta > 0$  we have

- (i)  $g_q(\beta, \xi) = g_q\left(\frac{\xi}{(\sqrt{\beta})^{2-q}}\right)$ ;
- (ii)  $g_q(\xi) = \frac{1}{C_q} Y_q(C_q^{1-q} \xi)$ .

**Proof.** The proof straightforwardly follows from the properties of the operator  $F_q$  indicated in proposition 2.1. □

These two formulas yield

$$F_q[G_q(\beta, x)](\xi) = \frac{1}{C_q} Y_q\left(\left(\frac{C_q}{\sqrt{\beta}}\right)^{1-q} \frac{\xi}{\sqrt{\beta}}\right). \tag{11}$$

Moreover,  $g_q(\beta, 0) = 1$ , which implies  $g_q(0) = 1$  and  $Y_q(0) = C_q$ . Thus, in order to know the properties of the  $q$ -Fourier transform of  $q$ -Gaussians it suffices to study  $Y_q(\xi)$ .

**Theorem 3.2.** Let  $1 \leq q < 3$  and  $q_n, n \in \mathbb{N}_q$ , are defined in equation (8). Then  $Y_{q_n}(\xi)$  satisfies the following homogeneous functional differential equation:

$$2\sqrt{q_n} \frac{dY_{q_n}(\xi)}{d\xi} + \xi Y_{q_{n-2}}(\sqrt{q_n} \xi) = 0, \quad \xi \in \mathbb{R}. \tag{12}$$

**Proof.** Differentiating  $Y_q(\xi) = \int_{-K}^K e_q^{-x^2+ix\xi} dx$  with respect to  $\xi$ , we have

$$\frac{dY_q(\xi)}{d\xi} = i \int_{-K}^K x (e_q^{-x^2+ix\xi})^q dx.$$

Further, integrating by parts,

$$\frac{dY_q(\xi)}{d\xi} = \frac{-i}{2} \int_{-K}^K d(e_q^{-x^2+ix\xi}) - \frac{\xi}{2} \int_{-K}^K (e_q^{-x^2+ix\xi})^q dx. \tag{13}$$

Obviously, the first integral vanishes if  $q \geq 1$ . Further, using  $(e_q^y)^q = e_{2-1/q}^{qy}$ , which is valid for any  $q < 3$  (see [6]), the second integral can be represented in the form

$$\int_{-K}^K (e_q^{-x^2+ix\xi})^q dx = \frac{1}{\sqrt{q}} \int_{-K}^K e_{2-1/q}^{-x^2+ix\sqrt{q}\xi} dx = \frac{1}{\sqrt{q}} Y_{2-1/q}(\sqrt{q}\xi). \tag{14}$$

Hence, for  $q \geq 1$  the function  $F_q[e_q^{-x^2}](\xi)$  satisfies the functional differential equation

$$2\sqrt{q} \frac{dY_q(\xi)}{d\xi} + \xi Y_{2-1/q}(\sqrt{q}\xi) = 0. \tag{15}$$

Setting  $q = q_n, n \in \mathbb{N}_q$ , and taking into account the relation  $2 - 1/q_n = q_{n-2}$ , we obtain equation (12). □

**Theorem 3.3.** Let  $0 < q < 1$  and  $q \neq l/(l + 1), l = 1, 2, \dots$ . Then  $Y_{q_n}(\xi)$  satisfies the following inhomogeneous functional differential equation:

$$2\sqrt{q_n} \frac{dY_{q_n}(\xi)}{d\xi} + \xi Y_{q_{n-2}}(\sqrt{q_n}\xi) = r_{q_n} \xi^{\frac{1}{1-q_n}}, \quad \xi \in \mathbb{R}, \tag{16}$$

where

$$r_{q_n} = 2\sqrt{q_n} \sin \frac{\pi}{2(1-q_n)} (1-q_n)^{\frac{1}{2(1-q_n)}}. \tag{17}$$

**Proof.** Assume that  $0 < q < 1$  and  $q \neq \frac{l}{l+1}, l = 1, 2, \dots$ . In this case the first integral on the right-hand side of (13) does not vanish, and takes the form

$$\int_{-K}^K d(e_q^{-x^2+ix\xi}) = e_q^{-K^2+iK\xi} - e_q^{-K^2-iK\xi} = 2i \operatorname{Im} e_q^{-K^2+iK\xi}.$$

Since  $\operatorname{supp} e_q^{-x^2} = [-K, K]$ , one has  $e_q^{-K^2} = 0$ . Hence,

$$e_q^{-K^2+iK\xi} = 0 \otimes_q e_q^{iK\xi} = [i(1-q)K\xi]^{\frac{1}{1-q}}.$$

Further, taking into account that  $K = 1/\sqrt{1-q}$ , one obtains

$$\operatorname{Im}[i(1-q)K\xi]^{\frac{1}{1-q}} = (1-q)^{\frac{1}{2(1-q)}} \sin \frac{\pi}{2(1-q)} \xi^{\frac{1}{1-q}}.$$

Note that the second integral on the right-hand side of equation (13) is the same as in the case of  $1 < q < 3$ . Consequently,  $F_q[e_q^{-x^2}](\xi)$  satisfies the functional differential equation

$$2\sqrt{q} \frac{dY_q(\xi)}{d\xi} + \xi Y_{2-1/q}(\sqrt{q}\xi) = r_q \xi^{\frac{1}{1-q}}, \tag{18}$$

where

$$r_q = 2\sqrt{q}(1-q)^{\frac{1}{2(1-q)}} \sin \frac{\pi}{2(1-q)}.$$

Again, setting  $q = q_n, n \in \mathbb{N}_q$ , we arrive at the functional differential equation (16). □

**Remark 3.4.**

- (i) If  $q = 0$ , then it is readily seen that  $Y_0(\xi) = F_0[e_0^{-x^2}](\xi) = \int_{-1}^1 (1-x^2+ix\xi) dx = 4/3$  for all  $\xi \in \mathbb{R}$ . Obviously, such  $Y_0(\xi)$  cannot be a  $q'$ -Gaussian for any  $q'$ .
- (ii) We will show later that a  $q$ -Fourier image of any  $q$ -Gaussian with  $q < 0$  cannot be a function of the form  $ae_q^{-\beta\xi^2}$ , for any  $q' \in (-\infty, 3)$  (see theorem 3.16).

Let us now consider the cases  $q = \ell/(\ell + 1), \ell = 1, 2, \dots$ , excluded from theorem 3.3. For these values of  $q$  we have  $K = \sqrt{\ell + 1}$  and

$$Y_q(\xi) = F_q[e_q^{-x^2}](\xi) = \int_{-\sqrt{\ell+1}}^{\sqrt{\ell+1}} \left(1 - \frac{1}{\ell+1}x^2 + \frac{1}{\ell+1}ix\xi\right)^{\ell+1} dx.$$

The latter is a polynomial of order  $\ell$  if  $\ell$  is even, and of order  $\ell + 1$  if  $\ell$  is odd<sup>3</sup>. In order to reflect this fact we use the conventional notation  $P_{\ell+1}(\xi) = Y_{\ell/(\ell+1)}(\xi)$  indicating the dependence on  $\ell$ . Further, obviously  $2 - \frac{1}{q} = \frac{\ell-1}{\ell}$ . Consequently,

$$Y_{2-1/q}(\xi) = \int_{-\sqrt{\ell}}^{\sqrt{\ell}} \left(1 - \frac{1}{\ell}x^2 + \frac{1}{\ell}ix\xi\right)^{\ell} dx = P_{\ell}(\xi),$$

<sup>3</sup>  $P_{\ell}$  does not contain odd-order terms.

and we note that  $P_\ell(\xi)$  is a polynomial of order  $\ell$  if  $\ell$  is even, and of order  $\ell - 1$  if  $\ell$  is odd. Moreover,  $P_\ell(\xi)$  is a symmetric function of  $\xi$  and  $P_\ell(0) = C_{\frac{\ell-1}{\ell}} > 0$ . Let  $\rho$  be a root of  $P_\ell(\xi)$  closest to the origin. We will consider  $P_\ell(\xi)$  only on its positivity interval  $(-\rho, \rho)$ .

**Theorem 3.5.** *Let  $q = \frac{2m-1}{2m}$ ,  $m = 1, 2, \dots$ . Then  $Y_q(\xi)$  satisfies equation (12).*

**Proof.** Assume  $\ell + 1 = 2m$ ,  $m = 1, 2, \dots$ . In this case  $Y_q(\xi) = P_{2m}(\xi)$  is a polynomial of order  $2m$  and  $Y_{2-1/q}(\xi) = P_{2m-1}(\xi)$  is a polynomial of order  $2m - 2$ . Moreover, it is easy to check  $r_q = 0$ . Thus,  $Y_q(\xi)$  satisfies the consistent equation

$$2\sqrt{q} \frac{dY_q(\xi)}{d\xi} + \xi Y_{2-1/q}(\sqrt{q}\xi) = 0, \quad \xi \in \mathbb{R}. \tag{19}$$

□

**Theorem 3.6.** *Let  $q = \frac{2m}{2m+1}$ ,  $m = 1, 2, \dots$ . Then  $Y_q(\xi)$  satisfies neither equation (12) nor (16).*

**Proof.** Let  $\ell = 2m$ ,  $m = 1, 2, \dots$ . Then  $Y_q(\xi) = P_{2m+1}(\xi)$  is a polynomial of order  $2m$ , and so is  $Y_{2-1/q}(\xi) = P_{2m}(\xi)$ . Assume that  $Y_q(\xi)$  satisfies equation (12), which in this particular case takes the form

$$2\sqrt{q} \frac{dY_q(\xi)}{d\xi} + \xi P_{2m}(\xi) = 0. \tag{20}$$

Equation (20) is clearly inconsistent, since the derivative of a polynomial of order  $2m$  cannot be a polynomial of order  $2m + 1$ . Analogously,  $Y_q(\xi)$  cannot satisfy equation (16) either. Indeed, if  $Y_q(\xi)$  solves equation (16), then in this particular case the equation would read

$$2\sqrt{q} \frac{dY_q(\xi)}{d\xi} + \xi P_{2m}(\xi) = \frac{(-1)^m}{(2m-1)^{m-\frac{1}{2}}} \xi^{2m+1}. \tag{21}$$

Equation (21) is inconsistent, since the term of the highest order on the left-hand side is  $\frac{2(-1)^m}{(2m+1)(2m)^{m-1/2}} \xi^{2m+1}$ , which is clearly distinct from the term of the highest order on the right-hand side. □

**Remark 3.7.** Equations (12) and (16) can be easily generalized for the  $q$ -Fourier transform of  $q$ -Gaussians with nonzero means. Namely, let  $\mu \neq 0$  be a real number, and

$$Y_{\mu,q}(\xi) = \int_{\mu-K}^{\mu+K} e_q^{-(x-\mu)^2+ix\xi} dx.$$

Then the associated functional differential equation for  $Y_{\mu,q_n}$  with  $q_n \in (0, 3)$  takes the form

$$2\sqrt{q_n} \frac{dY_{\mu,q_n}(\xi)}{d\xi} + \xi Y_{\frac{\mu}{q_n}, q_{n-2}}(\sqrt{q_n}\xi) - 2i\mu\sqrt{q_n} Y_{\mu,q_n}(\xi) = I_{(0,1)}(q_n) r_{q_n} \xi^{\frac{1}{1-q_n}}. \tag{22}$$

### 3.2. Is the $q$ -Fourier transform of a $q$ -Gaussian a $q'$ -Gaussian?

In this section we discuss a question important from applications point of view. Namely, we prove when the  $q$ -Fourier transform of a  $q$ -Gaussian is a  $q'$ -Gaussian with some index  $q'$  in  $(-\infty, 3)$ . With this aim we introduce the set of functions

$$\mathcal{G} = \bigcup_{q < 3} \mathcal{G}_q, \quad \text{where } \mathcal{G}_q = \{f : f(x) = ae_q^{-\beta x^2}, a > 0, \beta > 0\}. \tag{23}$$

It follows from relationship (11) that if the  $q$ -Fourier transform  $F_q[G_q(\beta, x)](\xi)$  of a  $q$ -Gaussian is a  $q'$ -Gaussian with some  $q' \in (-\infty, 3)$ , then  $Y_q(\xi)$  must belong to  $\mathcal{G}$ . Therefore,



we will study the existence of a solution of functional differential equations (12) and (16) in the set  $\mathcal{G}$ .

**Theorem 3.8.** *Let  $1 \leq q < 3$  and  $q_n, n \in \mathbb{N}_q$ , be the sequence defined in (8). Then the functional differential equation*

$$2\sqrt{q_n} \frac{dY_{q_n}(\xi)}{d\xi} + \xi Y_{q_{n-2}}(\sqrt{q_n}\xi) = 0, \quad \xi \in \mathbb{R}, \tag{24}$$

has a unique solution  $Y_{q_n}(\xi) \in \mathcal{G}$  satisfying the condition

$$Y_{q_n}(0) = C_{q_n}. \tag{25}$$

This solution is specifically

$$Y_{q_n}(\xi) = C_{q_n} e_{q_{n+1}}^{-\frac{3-q_n}{8}\xi^2}. \tag{26}$$

**Proof.** *Existence.* It follows immediately from representation (26) that  $Y_{q_n}(0) = C_{q_n}$ . Furthermore,

$$\frac{dY_{q_n}(\xi)}{d\xi} = -\frac{1}{4}(3-q_n)C_{q_n}\xi \left( e_{q_{n+1}}^{-\frac{3-q_n}{8}\xi^2} \right)^{q_{n+1}}, \tag{27}$$

$$Y_{q_{n-2}}(\sqrt{q_n}\xi) = C_{q_{n-2}} e_{q_{n-1}}^{-q_n \frac{3-q_{n-2}}{8}\xi^2}. \tag{28}$$

Due to the equation  $(e_q^y)^q = e_{2-1/q}^{qy}$  and part (i) of proposition 2.6, expression (27) can be rewritten as

$$\frac{dY_{q_n}(\xi)}{d\xi} = -\frac{1}{4}(3-q_n)C_{q_n}\xi e_{q_{n-1}}^{-q_n \frac{3-q_{n-2}}{8}\xi^2}. \tag{29}$$

Substituting (28) and (29) into (24), we obtain

$$\left( -\sqrt{q_n} C_{q_n} \frac{3-q_n}{2} + C_{q_{n-2}} \right) e_{q_{n-1}}^{-\frac{q_n(3-q_n)}{8}\xi^2} = 0. \tag{30}$$

Now taking into account part (ii) of proposition 2.6 we conclude that  $Y_{q_n}(\xi)$  in (26) satisfies (24).

*Uniqueness.* We note that  $|\cos_q(x)| \leq 1$  for real  $x$ , if  $q > 1$  (see [6]). This fact and corollary 2.3 imply the following estimate:

$$|Y_q(\xi)| = \left| \int_{-\infty}^{\infty} e_q^{-x^2+ix\xi} dx \right| \leq \int_{-\infty}^{\infty} e_q^{-x^2} dx = C_q. \tag{31}$$

Assume that there are two solutions to problem (24)–(25), i.e.  $Y_{q_n}$  and  $\tilde{Y}_{q_n}$ . Then their difference  $Z_{q_n}(\xi) = Y_{q_n}(\xi) - \tilde{Y}_{q_n}(\xi)$  also satisfies equation (24), and the condition  $Z_{q_n}(0) = 0$ . Now estimate (31) yields  $Z_{q_n} \equiv 0$ , which, in turn, implies  $Y_{q_n} \equiv \tilde{Y}_{q_n}$ .  $\square$

**Corollary 3.9.** *Let  $q_n \geq 1$ . Then*

$$F_{q_n}[G_{q_n}](\xi) = e_{q_{n+1}}^{-\frac{3-q_n}{8\beta^{2-q_n} C_{q_n}^{2(q_n-1)}}\xi^2}. \tag{32}$$

**Remark 3.10.** Representation (32) was obtained in [6] by the contour integration technique. The formula in (7) corresponds to the particular case  $n = 0$  of (32).

**Remark 3.11.** If  $q = 1$ , then the Cauchy problem (24)–(25) reads

$$2 \frac{dY_1(\xi)}{d\xi} + \xi Y_1(\xi) = 0, \quad Y_1(0) = \sqrt{\pi},$$

and its unique solution is  $Y_1(\xi) = \sqrt{\pi} e^{-\xi^2/4}$ . Besides corollary 3.9 we obtain

$$F \left[ \frac{\sqrt{\beta}}{\sqrt{\pi}} e^{-\beta x^2} \right] = e^{-\frac{1}{4\beta} \xi^2}.$$

The density of the standard normal distribution corresponds to  $\beta = 1/2$ , giving the characteristic function of the classic Gaussian.

**Theorem 3.12.** Let  $q_n, n \in \mathbb{N}_q$ , be a sequence defined in (8) with  $q < 1$ . Suppose that  $q_n \neq m/(m + 1), m = 1, 2, \dots$ . Then the functional differential equation

$$2\sqrt{q_n} \frac{dY_{q_n}(\xi)}{d\xi} + \xi Y_{q_{n-2}}(\sqrt{q_n}\xi) = r_{q_n} \xi^{1-q_n}, \quad \xi \in \mathbb{R}, \quad (33)$$

has no solution in  $\mathcal{G}$ .

**Proof.** We recall that if  $q < 1$  and  $n \in \mathbb{N}_q$ , then  $q_n < 1$ . Further, we note that a function with compact support cannot solve equation (33). Since each function in  $\mathcal{G}_q$  for  $q < 1$  has a compact support, the solution of equation (33) cannot belong to  $\mathcal{G}_q, q < 1$ . Now assume that  $Y_{q_n}(\xi) \in \mathcal{G}_{q'}$  and  $Y_{q_{n-2}}(\xi) \in \mathcal{G}_{q''}$ , with  $q' > 1$  or  $q'' > 1$  (the reader can easily verify that  $q' \neq 1$  and  $q'' \neq 1$ ). In accordance with definition (23),

$$Y_{q_n}(\xi) = a e_{q'}^{-b\xi^2} \quad \text{and} \quad Y_{q_{n-2}}(\xi) = A e_{q''}^{-B\xi^2},$$

where  $a, b, A, B$  are some real positive numbers depending on  $q_n$ . Then, for equation (33) to be consistent,

$$\frac{2}{1-q'} - 1 = \frac{1}{1-q_n} \quad \text{or} \quad \frac{2}{1-q''} + 1 = \frac{1}{1-q_n}.$$

Solving these equations for  $q'$  and  $q''$ , one obtains  $q' = \frac{q_n}{2-q_n}$  and  $q'' = \frac{3q_n-2}{q_n}$ . It follows  $\max(q', q'') < 1$ , since  $q_n < 1$ . This contradicts the assumption that  $q' > 1$ , or  $q'' > 1$ .  $\square$

Next, we consider the cases  $q = \frac{1}{2}, \frac{2}{3}, \dots, \frac{m}{m+1}, \dots$ , excluded from theorem 3.12. Direct computations show that in two specific cases, namely  $q = 1/2$  and  $q = 2/3, Y(q, \xi) \in \mathcal{G}_0$  considered on the positivity intervals. Indeed,

$$Y_{\frac{1}{2}}(\xi) = F_{\frac{1}{2}} \left[ e_{\frac{1}{2}}^{-x^2} \right](\xi) = \frac{16\sqrt{2}}{15} \left( 1 - \frac{5}{16} \xi^2 \right), \quad (34)$$

which is nonnegative for  $\xi \in [-4/\sqrt{5}, 4/\sqrt{5}]$ . Therefore, on this interval we can associate it with an element of  $\mathcal{G}_0$ , writing  $Y(1/2, \xi) = \frac{16\sqrt{2}}{15} e_0^{-(5/16)\xi^2} \in \mathcal{G}_0$ . Similarly,

$$Y_{\frac{2}{3}}(\xi) = F_{\frac{2}{3}} \left[ e_{\frac{2}{3}}^{-x^2} \right](\xi) = \frac{32\sqrt{3}}{35} \left( 1 - \frac{7}{24} \xi^2 \right) \in \mathcal{G}_0, \quad (35)$$

on the positivity interval  $(-2\sqrt{\frac{6}{7}}, 2\sqrt{\frac{6}{7}})$ .

However,  $Y(q, \xi)$  does not belong to  $\mathcal{G}$  for any other value of  $q = 3/4, 4/5, \dots$ . In order to show this first we derive an explicit form for  $P_{m+1}(\xi) = Y_{m/(m+1)}(\xi)$ . Recall that  $P_{m+1}(\xi)$  is a polynomial of order  $m + 1$  if  $m + 1$  is even. Otherwise, it is a polynomial of order  $m$ .

**Theorem 3.13.** Let  $q = m/(m + 1)$ ,  $m = 1, 2, \dots$ . Then  $Y_q(\xi) = P_{m+1}(\xi)$  has the representation

$$P_{m+1}(\xi) = \sum_{k=0}^{\lfloor \frac{m+1}{2} \rfloor} (-1)^k (m + 12k) (m + 1)^{-k+\frac{1}{2}} B\left(k + \frac{1}{2}, m - 2k + 2\right) \xi^{2k}, \tag{36}$$

where  $[x]$  is the integer part of  $x$ , and  $B(a, b)$  is Euler’s beta-function.

**Proof.** Recall that if  $q = \frac{m}{m+1}$ , then  $Y_q(\xi)$  has the form

$$Y_q(\xi) = P_{m+1}(\xi) = \int_{-\sqrt{m+1}}^{\sqrt{m+1}} \left(1 - \frac{1}{m+1}x^2 + \frac{1}{m+1}ix\xi\right)^{m+1} dx.$$

We have

$$P_{m+1}(\xi) = \sum_{k=0}^{m+1} (m + 1k) D_k(m) \frac{(i\xi)^k}{(m + 1)^k},$$

where

$$D_k(m) = \int_{-\sqrt{m+1}}^{\sqrt{m+1}} \left(1 - \frac{1}{m+1}x^2\right)^{m-k+1} x^k dx.$$

It is not hard to verify that  $D_k(m) = 0$  if  $k$  is odd and

$$D_{2k}(m) = (m + 1)^{k+1/2} B(k + 1/2, m - 2k + 2)$$

for  $k = 0, \dots, \lfloor \frac{m+1}{2} \rfloor$ , which implies representation (36). □

**Theorem 3.14.** Let  $q = m/(m + 1)$ ,  $m = 3, 4, \dots$ . Then  $Y_q(\xi) \notin \mathcal{G}$ .

**Proof.** It follows from representation (36) that the polynomial  $Y_q(\xi) = P_{m+1}(\xi)$ , with the first three (nonzero) terms indicated, reads

$$\begin{aligned} Y_q(\xi) &= D_0(m) \left[ 1 - (m + 1)^2 \frac{B(\frac{3}{2}, m)}{B(\frac{1}{2}, m + 2)} \xi^2 + \frac{m(m + 1)^3}{2} \frac{B(\frac{5}{2}, m - 2)}{B(\frac{1}{2}, m + 2)} \xi^4 + \dots \right] \\ &= D_0(m) \left[ 1 - \frac{2m + 3}{8(m + 1)} \xi^2 + \frac{(2m + 3)(2m + 1)}{8(m + 1)^2} \xi^4 + \dots \right], \end{aligned} \tag{37}$$

where

$$D_0(m) = C_{\frac{m}{m+1}} = \sqrt{m + 1} B\left(\frac{1}{2}, m + 2\right) = \frac{\sqrt{m + 1}(m + 1)!2^{m+2}}{(2m + 3)!!}.$$

Now assume that  $Y_q(\xi) \in \mathcal{G}_{q'}$  for some  $q' < 3$ . Then  $1/(1 - q') = (m + 1)/2$ , or  $q' = (m - 1)/(m + 1)$ . Therefore,

$$Y_q(\xi) = D_0(m)[1 - \beta(m)\xi^2]^{1+\frac{m+1}{2}},$$

where  $\beta(m) > 0$  and  $|\xi| \leq 1/\sqrt{\beta(m)}$ . Applying the binomial formula and indicating the first three terms, one has

$$Y_q(\xi) = D_0(m) \left[ 1 - \frac{(m + 1)\beta(m)}{2} \xi^2 + \frac{(m^2 - 1)[\beta(m)]^2}{8} \xi^4 + \dots \right]. \tag{38}$$

Comparing the second and third terms in (37) and (38), one obtains contradictory relations

$$\beta(m) = \frac{2m + 3}{4(m + 1)^2}$$

and

$$[\beta(m)]^2 = \frac{(3m+3)(2m+1)}{(m-1)(m+1)^3} \neq \frac{(2m+3)^2}{16(m+1)^4} = [\beta(m)]^2, \quad m = 3, 4, \dots,$$

which completes the proof. □

**Remark 3.15.** Formula (36) for  $q = 1/2$  and  $q = 2/3$  gives

$$Y_{\frac{1}{2}}(\xi) = \frac{16\sqrt{2}}{15} \left(1 - \frac{5}{16}\xi^2\right) = \frac{16\sqrt{2}}{15} e_0^{-(5/16)\xi^2}, \quad \xi \in \left[-\frac{4\sqrt{5}}{5}, \frac{4\sqrt{5}}{5}\right],$$

and

$$Y_{\frac{2}{3}}(\xi) = \frac{32\sqrt{3}}{35} \left(1 - \frac{7}{24}\xi^2\right) = \frac{32\sqrt{3}}{35} e_0^{-\frac{7}{24}\xi^2}, \quad \xi \in \left[-2\sqrt{\frac{7}{6}}, 2\sqrt{\frac{7}{6}}\right].$$

which coincide with (34) and (35), respectively. Both functions belong to  $\mathcal{G}_0$ .

**Theorem 3.16.** Let  $q < 0$ . Then  $Y_q(\xi) \notin \mathcal{G}$ .

**Proof.** Repeating calculations used in proofs of theorems 3.2 and 3.3, it is not hard to verify that the derivative of  $Y_q(\xi)$  can be represented in the form

$$\frac{dY_q(\xi)}{d\xi} = -\frac{\xi}{2} \int_{-K}^K (e_q^{-x^2+ix\xi})^q dx + R_q \xi^{\frac{1}{1-q}}, \tag{39}$$

where

$$R_q = (1-q)^{\frac{1}{2(1-q)}} \sin \frac{\pi}{2(1-q)}. \tag{40}$$

We note that the condition  $q < 0$  implies two statements important for further proof, namely  $0 < \frac{1}{1-q} < 1$ , and  $R_q \neq 0$ . Now assume that  $Y_q \in \mathcal{G}_{q'}$  with some  $q' \in (-\infty, 3)$ . In other words, there are positive numbers  $a$  and  $b$ , such that  $Y_q(\xi) = a e_{q'}^{-b\xi^2}$ . Taking the first derivative of the latter, equating it to the right-hand side of (39), and dividing both sides by  $\xi^{\frac{1}{1-q}}$  ( $\xi \neq 0$ ), we obtain

$$\xi^{-\frac{q}{1-q}} \left[ \frac{1}{2} \int_{-K}^K (e_q^{-x^2+ix\xi})^q dx - 2ab(e_{q'}^{-b\xi^2})^{q'} \right] = R_q, \tag{41}$$

which must be valid for all  $\xi \in \mathbb{R}$ . However, the left-hand side becomes arbitrarily small for small  $\xi$ , since  $-\frac{q}{1-q} > 0$  and the expression in brackets has a finite limit at 0, while the right-hand side is nonzero constant. This contradiction completes the proof. □

#### 4. Some applications to the porous medium equation

In this section we discuss some applications of the  $q$ -Fourier transform  $F_q$  to nonlinear models of partial differential equations. First we verify that the theorems proved in section 3 imply that  $F_q$  transfers a  $q$ -Gaussian into a  $q_1$ -Gaussian if  $q \geq 1$ ,  $q_1 = (1+q)/(3-q)$ . Moreover, as shown in [11], the operator  $F_q : G_q \rightarrow G_{q_1}$ , for  $q > 1$  is invertible. These two facts have been essentially used in [6, 10] for the proof of  $q$ -versions of the central limit theorem. Another application of  $F_q$ , as sketched hereunder, shows that it can be used for establishing a relation between the porous medium equation and a nonlinear ordinary differential equation (ODE) similar to the usual Fourier transform.

The classic Fourier transform reduces the Cauchy problem for linear partial differential equations of the form  $u_t(t, x) = A(D_x)u(t, x)$ ,  $t > 0$ ,  $x \in \mathbb{R}^n$ ,  $u(0, x) = \varphi(x)$ ,  $x \in \mathbb{R}^n$ , where

$D_x = (D_1, \dots, D_n)$ ,  $D_j = -i \frac{\partial}{\partial x_j}$ ,  $j = 1, \dots, n$ , and  $A(D_x)$  is an elliptic differential operator, to an associated linear ODE with the parameter  $\xi \in R^n$ . In the particular case of  $n = 1$  and  $A(D_x) = \frac{d^2}{dx^2}$  for the Fourier image  $\hat{u}(t, \xi)$  of a solution  $u(t, x)$ , we have a dual differential equation

$$\hat{u}'_t(t, \xi) = -\xi^2 \hat{u}(t, x), \quad \hat{u}(0, \xi) = \hat{\varphi}(\xi), \tag{42}$$

where  $\xi \in R^1$  is a parameter. This instance corresponds to the Fokker–Planck equation without drift [15] for some classes of nonlinear differential equations.

We now demonstrate the similar role of  $F_q$  in the celebrated *porous medium equation* in the superdiffusion regime ubiquitously found in physical phenomena [17–21] (and references therein)<sup>4</sup>. Consider the following nonlinear diffusion equation with a singular diffusion coefficient:

$$\frac{\partial U}{\partial t} = (U^{1-q} U_x)_x, \quad t > 0, \quad x \in R^1, \quad q > 1. \tag{43}$$

We look for a solution in the similarity set  $G_q^* = \{U(t, x) : U(t, x) = t^a G_q(\beta; t^b x), a = a(q), b = b(q) \in R^1, \beta = \beta(q) > 0\}$ , where  $a$  and  $\beta$  do not depend on  $t$  and  $x$ .

**Proposition 4.1.** *Suppose  $U(t, x) \in G_q^*$  is a solution to equation (43). Then its  $q$ -Fourier transform  $\hat{U}_q(t, \xi) = F_q[U(t, x)](\xi)$  satisfies the following nonlinear ordinary differential equation with the parameter  $\xi$ :*

$$(\hat{U}_q)'_t = -\frac{B(\beta, q)\xi^2}{t^{\frac{q-1}{3-q}}} (\hat{U}_q)^{q_1}, \quad t > 0, \tag{44}$$

where  $B(\beta, q) = \frac{2-q}{4\beta^{2-q} C_q^{q-1}}$  and  $q_1 = \frac{1+q}{3-q}$ .

**Proof.** Let  $U \in G_q^*$  be a solution to (43), i.e. for some  $a = a(q)$  and  $\beta = \beta(q)$  it has the representation  $U(t, x) = t^a G_q(\beta; t^a x)$ . Then it follows from proposition 2.1 that

$$\begin{aligned} \hat{U}_q(t, \xi) &= F_q[U(t, x)](\xi) \\ &= F_q[G_q(\beta; x)]\left(\frac{\xi}{t^{a(2-q)}}\right) = \frac{1}{C_q} Y_q\left(\left(\frac{\sqrt{\beta}}{C_q}\right)^{q-1} \frac{\xi}{\sqrt{\beta} t^{a(2-q)}}\right), \end{aligned}$$

where  $Y_q(\xi)$  is a solution to equation (24). Computing the derivative of  $\hat{U}_q(t, x)$  in variable  $t$ , taking into account that  $a = -1/(3 - q)$  (see, e.g., [21]), and using equation (24), we obtain

$$(\hat{U}_q)'_t = -\frac{2-q}{4\beta^{2-q} C_q^{2(q-1)}} \xi^2 (\hat{U}_q)^{q_1},$$

where  $q_1 = (1 + q)/(3 - q)$ . □

The inverse statement, given in the following formulation, is also true.

**Proposition 4.2.** *Suppose  $V(t, \xi)$ ,  $V(0, \xi) = 1$ , is a solution to ODE with the parameter  $\xi$*

$$V' = -\frac{B(\beta, q)\xi^2}{t^{\frac{q-1}{3-q}}} V^{q_1}, \quad t > 0, \tag{45}$$

where  $B(q, \beta)$  and  $q_1$  are as in proposition 4.1. Then its inverse  $q$ -Fourier transform  $U(t, x) = F_q^{-1}[V(t, \xi)](x)$  exists and satisfies equation (43).

<sup>4</sup> The monograph [21] contains different approaches to the solution of the porous medium equation.

**Proof.** By separation of variables of (45) one can verify that its solution

$$V(t, \xi) = e_{q_1}^{-\frac{3-q}{8\beta^{2-q}C_q^{q-1}}(\xi t^{\frac{2-q}{3-q}})}.$$

By theorem 0.6 of paper [11] the inverse  $q$ -Fourier transform for  $V(t, \xi)$  exists, and by virtue of propositions 2.1 and 2.4 it has the representation

$$U(t, x) = \frac{1}{t^{\frac{1}{3-q}}} G_q \left( \beta(q); \frac{x}{t^{\frac{1}{3-q}}} \right), \quad \text{where } \beta(q) = \frac{1}{[2(3-q)C_q^{\frac{1}{q-1}}]^{\frac{2}{3-q}}}. \quad (46)$$

The latter is a solution to equation (43); see [21]. □

Note that if the initial condition is given in the form  $U(0, x) = \delta(x)$  with the Dirac delta function, and  $q = 1$ , then we obtain (42) ( $\hat{\varphi}(\xi) \equiv 1$ ), in which  $\beta = 1/4$ ,  $B(\beta, 1) = 4\beta = 1$ .

In order to study price fluctuations in stock markets a stochastic process  $X_t = \frac{\ln S(t+t_0)}{\ln S(t_0)}$  representing log-returns was introduced in [16]. Here  $S(t)$  is the price at time  $t$ .  $X_t$  solves a stochastic differential equation  $dX_t = \tau dt + \sigma d\Omega_t$ , where  $\tau$  and  $\sigma$  are the drift and volatility coefficients respectively, and  $\Omega_t$  is a solution to the Itô stochastic differential equation

$$d\Omega_t = [P(\Omega_t)]^{\frac{1-q}{2}} dB_t, \quad t > t_0. \quad (47)$$

In this equation  $B_t$  is a Brownian motion, and  $P$  is a  $q$ -Gaussian distribution function. The corresponding Fokker–Planck-type equation in the case  $\tau = 0$ ,  $\sigma = 1$  reads

$$\frac{\partial V(x, t|x', t')}{\partial t} = ([V(x, t|x', t')]^{2-q})_{xx}, \quad (48)$$

which can easily be reduced to the form (43). From the financial applications point of view it is important to know the properties of the stochastic process  $X_t$ , since it can be considered as a  $q$ -alternative to the Brownian motion. One can effortlessly verify that if  $U(t, x)$  is a solution to equation (43) for  $t > 0$  with an initial condition  $U(0, x) = f(x)$ , then a solution  $V(t, x)$ ,  $t > t'$ , to the same equation (43) considered for  $t > t'$  with an initial condition  $V(t', x) = f(x)$  can be represented in the form  $V(t, x) = U(t - t', x)$ ,  $t > t'$ . It follows that  $X_t$  has stationary increments.

Concluding the discussion we note that solution (46) corresponds to the solution obtained from an ansatz [20] which is in accordance with the generalized central limit theorem presented in [6]. The method we have just presented for the model case can be implemented for other more general cases as well. For instance, the Fokker–Planck-type equation associated with a process  $X_t$  with constant drift  $\tau = \mu \neq 0$ , due to a term  $-2i\mu\sqrt{q_n} Y_{\mu, q_n}(\xi)$  in equation (22), has an additional drift term on the right-hand side of equation (48). We also note that with more routine calculations the method can be extended to the case of time-dependent drift and diffusion coefficients. We intend to present all the routine calculations in the general case of linear external forces in a separate paper.

### 5. Conclusion

Summarizing, we have the following general picture for the  $q$ -Fourier transform of  $q$ -Gaussians.

- (1) The case  $1 \leq q < 3$ :
  - (1a) the  $q$ -Fourier transform acts as  $F_q : \mathcal{G}_q \rightarrow \mathcal{G}_{q'}$ ;
  - (1b) the relation between  $q$  and  $q'$  is given by  $q' = \frac{1+q}{3-q}$ .

- (2) The case  $q = \frac{1}{2}$  or  $q = \frac{2}{3}$  : in this case the operator acts as  $F_q : \mathcal{G}_q \rightarrow \mathcal{G}_0$ . Relationship (1b) is failed.
- (3) The case  $q < 1$ , but  $q \neq \frac{1}{2}, \frac{2}{3}$  : in this case (1a) is failed, in the sense that there is no  $q'$  such that the  $q$ -Fourier transform of a  $q$ -Gaussian would be a  $q'$ -Gaussian.

The lesson we have learnt from the above analysis is that the  $q$ -Fourier transform defined by formula (1) (or, the same, by formula (3)) is rich in content and applications if  $q \in [1, 3)$ . Its important applications in the case  $q > 1$  are given in [6] for the proof of the  $q$ -central limit theorem, and in [9] for the classification of  $(q, \alpha)$ -stable distributions. Another application of the operator  $F_q$  to the porous medium equation and related stochastic differential models is discussed in section 4 of the current paper. What concerns the case  $q < 1$ , the  $q$ -Fourier transform defined by formula (1) does not possess the properties valid in the case  $1 \leq q < 3$ . Therefore, the methods developed for  $q \geq 1$  are not applicable in the case  $q < 1$ . The study of applications of the  $q$ -Fourier transform (or its alternatively defined version) to problems mentioned above, including the  $q$ -central limit theorem, remains a challenging open question in the case  $q < 1$ .

### Acknowledgments

We acknowledge C Tsallis for several comments on the subjects mentioned in this paper. We also acknowledge thoughtful remarks by anonymous referees which significantly improved the text.

### References

- [1] Tsallis C 1988 *J. Stat. Phys.* **52** 479  
 Curado E M F and Tsallis C 1991 *J. Phys. A: Math. Gen.* **24** L69  
 Curado E M F and Tsallis C 1991 **24** 3187 (corrigendum)  
 Curado E M F and Tsallis C 1992 **25** 1019 (corrigendum)
- [2] Tsallis C 2009 *Introduction to Nonextensive Statistical Mechanics: Approaching a Complex World* (Berlin: Springer)
- [3] Gell-Mann M and Tsallis C 2004 *Nonextensive Entropy—Interdisciplinary Applications* (New York: Oxford University Press)
- [4] Abe S and Okamoto Y (eds) 2001 *Nonextensive Statistical Mechanics and Its Applications (Lecture Notes in Physics vol 560)* (Heidelberg: Springer)
- [5] Tsallis C, Mendes R S and Plastino A R 1998 *Physica A* **261** 534  
 Kaniadakis G, Lissia M and Rapisarda A (eds) 2002 *Non-extensive Thermodynamics and Physical Applications (Physica A vol 305)* (Amsterdam: Elsevier)  
 Swinney H L and Tsallis C (eds) 2004 *Anomalous Distributions, Nonlinear Dynamics and Nonextensivity (Physica D vol 193)* (Amsterdam: Elsevier)  
 Abe S, Herrmann H J, Quarati P, Rapisarda A and Tsallis C (eds) 2007 *Complexity, Metastability and Nonextensivity (AIP Conf. Proc. vol 965)* (New York: American Institute of Physics)
- [6] Umarov S, Tsallis C and Steinberg S 2008 *Milan J. Math.* **76** 307
- [7] Tsallis C 2005 *Milan J. Math.* **73** 145
- [8] Moyano L G, Tsallis C and Gell-Mann M 2006 *Europhys. Lett.* **73** 813
- [9] Umarov S, Tsallis C, Gell-Mann M and Steinberg S 2010 *J. Math. Phys.* (at press)  
 Umarov S, Tsallis C, Gell-Mann M and Steinberg S 2006  $q$ -generalization of symmetric alpha-stable distributions: Part I arXiv:cond-mat/06006038
- [10] Umarov S and Tsallis C 2007 On multivariate generalizations of the  $q$ -central limit theorem consistent with nonextensive statistical mechanics *Complexity, Metastability and Nonextensivity (AIP Conf. Proc. vol 965)* ed S Abe, H J Herrmann, P Quarati, A Rapisarda and C Tsallis (New York: American Institute of Physics)
- [11] Umarov S and Tsallis C 2008 *Phys. Lett. A* **372** 29  
 Umarov S, Tsallis C, Gell-Mann M and Steinberg S 2006  $q$ -generalization of symmetric alpha-stable distributions: Part II arXiv:cond-mat/06006040

- [12] Nivanen L, Le Mehaute A and Wang Q A 2003 *Rep. Math. Phys.* **52** 437
- [13] Borges E P 2004 *Physica A* **340** 95
- [14] Borges E P 1998 *J. Phys. A: Math. Gen.* **31** 5281–8
- [15] Risken H 1989 *The Fokker–Planck Equation—Methods of Solution and Applications* 2nd edn (Berlin: Springer)
- [16] Borland L 2002 *Phys. Rev. Lett.* **89** 098701
- [17] Carrillo J A and Toscani G 2000 *Indiana Univ. Math. J.* **49** 113
- [18] Otto F 2001 *Commun. Part. Differ. Equ.* **26** 101
- [19] Muskat M 1937 *The Flow of Homogeneous Fluids Through Porous Media* (New York: McGraw-Hill)
- [20] Tsallis C and Bukman D J 1996 *Phys. Rev. E* **54** R2197
- [21] Vázquez J L 2007 *The Porous Medium Equation: Mathematical Theory* (*Oxford Mathematical Monographs*) (New York: Oxford University Press)